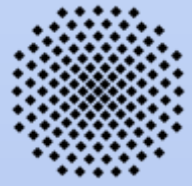


**Universität  
Stuttgart**

# thermodynamics of colloidal particles

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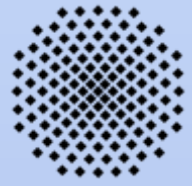
Thomas Speck  
Valentin Blickle  
Laurent Helden  
Udo Seifert  
Clemens Bechinger



# contents

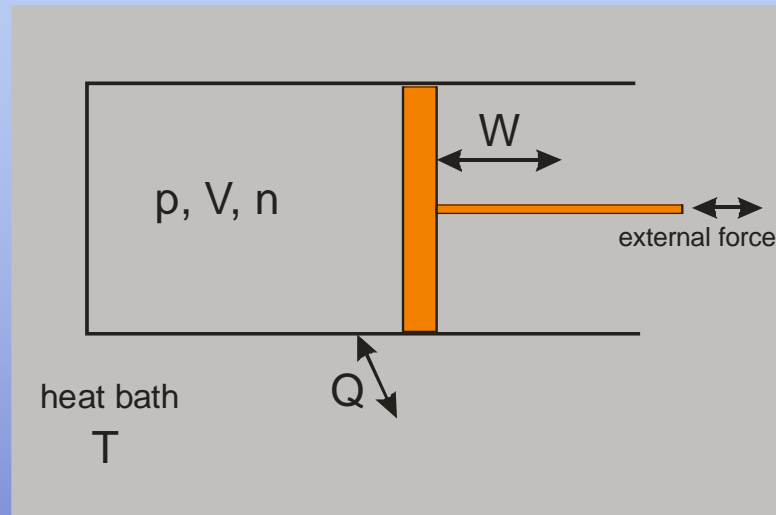
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- mesoscopic thermodynamics
- TIRM experiment
- energy conservation
- work distribution / fluctuation theorems



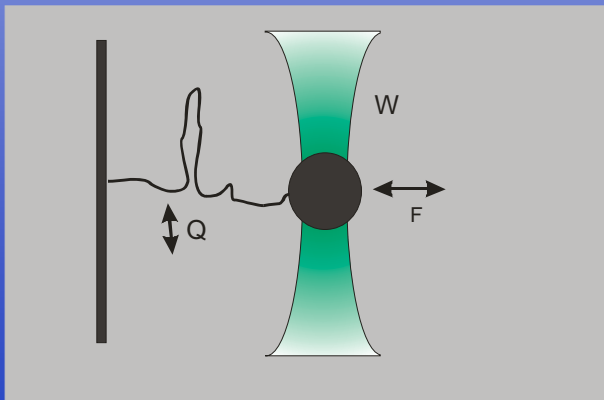
# thermodynamics

macroscopic:

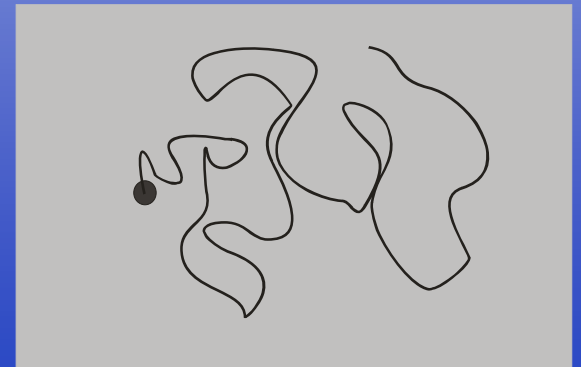


ideal gas

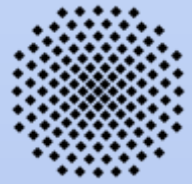
mesoscopic:



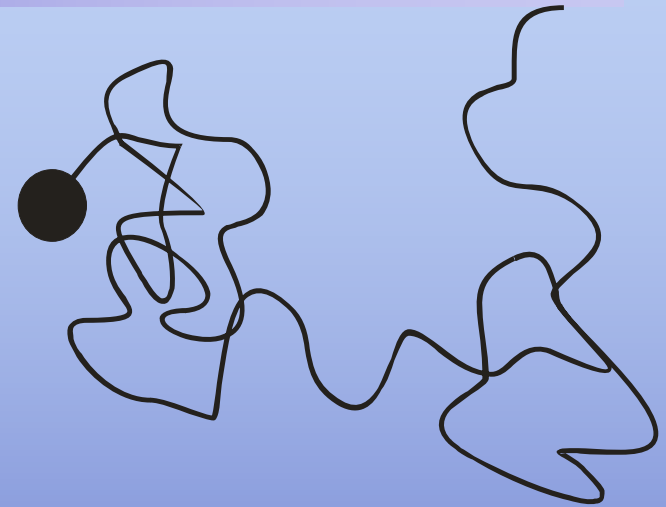
- suspended colloidal particle
- protein pulling
- Brownian motion
- fluctuations



# brownian motion



Brownian particle: random walk



Langevin equation

$$\cancel{m \frac{d^2 z}{dt^2}} + \gamma \frac{dz}{dt} = -\frac{\partial V}{\partial z} + \xi(t)$$

overdamped system

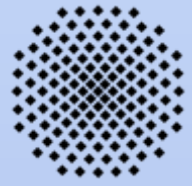
friction coefficient

external potential

stochastic force:

$$\langle \xi(t) \rangle = 0$$

$$\langle \xi(t) \cdot \xi(t') \rangle = 2\gamma T \delta(t - t')$$



# energy balance: equilibrium

time-independent potential:  $\frac{dV}{dt} = 0$        $V$  only  $z$  dependent

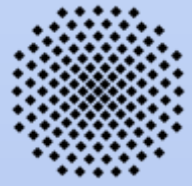
Langevin equation  $\gamma \frac{dz}{dt} = -\frac{\partial V}{\partial z} + \xi(t)$

$$0 = -\left[-\gamma \frac{dz}{dt} + \xi(t)\right] dz + \frac{\partial V}{\partial z} dz$$

$$dQ \equiv -\left[-\gamma \frac{dz}{dt} + \xi(t)\right] dz \quad \text{heat}$$

$$dV \equiv \frac{\partial V}{\partial z} dz \quad \text{potential difference}$$

$$0 = dQ + dU \quad \text{energy balance}$$



# energy balance: time dep. potentials

$V[z, \lambda(t)]$   $\lambda(t)$  control parameter

e.g.:  $\lambda(t) = \sin(\omega t)$

no total differential  
add:  $\frac{\partial V(z, \lambda(t))}{\partial \lambda} d\lambda$

$$0 = -\left[-\gamma \frac{dz}{dt} + \xi(t)\right] dz + \frac{\partial V(z, \lambda(t))}{\partial z} dz$$

$$\underbrace{\frac{\partial V(z, \lambda(t))}{\partial \lambda} d\lambda}_{dW} = \underbrace{-\left[-\gamma \frac{dz}{dt} + \xi(t)\right] dz}_{dQ} + \underbrace{\frac{\partial V(z, \lambda(t))}{\partial z} dz + \frac{\partial V(z, \lambda(t))}{\partial \lambda} d\lambda}_{dV}$$

work

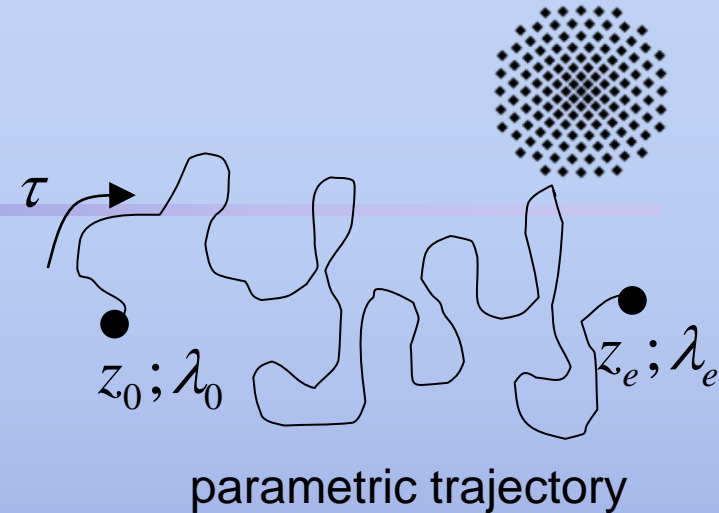
heat

potential difference

microscopic energy balance

# trajectory picture

$$dQ = -\left[-\gamma \frac{dz}{dt} + \xi(t)\right] dz = -\frac{\partial V}{\partial z} dz$$



Langevin equation

$$\gamma \frac{dz}{dt} = -\frac{\partial V}{\partial z} + \xi$$

integration along trajectory

$$Q = \int_{z_0, \lambda_0}^{z_e, \lambda_e} -\frac{\partial V}{\partial z} dz = \int_{t_0}^{t_e} -\frac{\partial V}{\partial z} \dot{z} d\tau$$

$$W = \int_{t_0}^{t_e} \frac{\partial V}{\partial \lambda} \dot{\lambda} d\tau$$

$$\Delta V = V(t_e) - V(t_0)$$

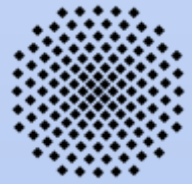
trajectory



determination of Q, W, and V

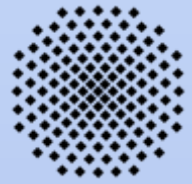
# experiment

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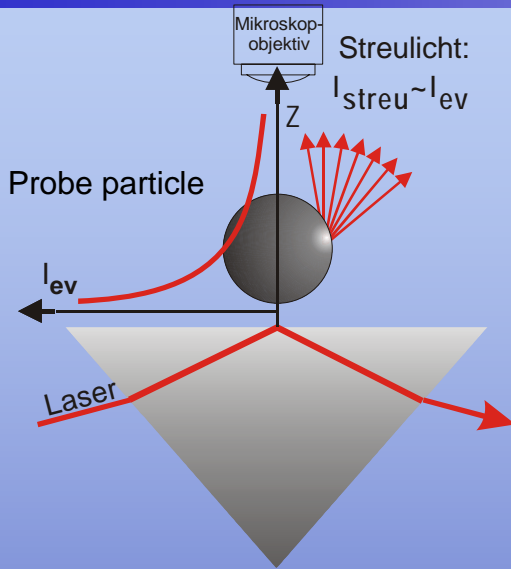


- exact measurement of the trajectory
- externally controllable force to drive particle into non equilibrium

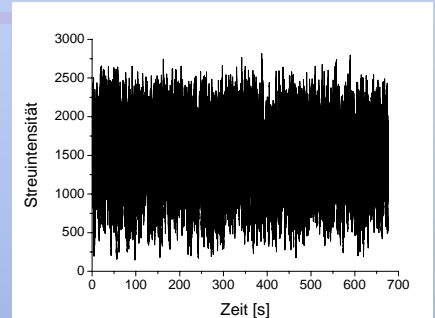




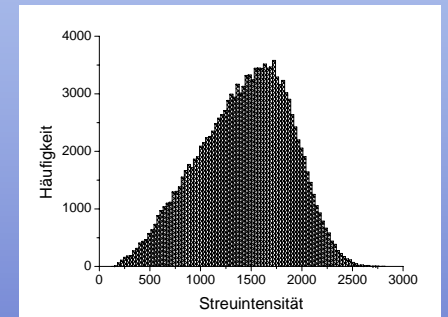
# TIRM Technique



Scattering signal

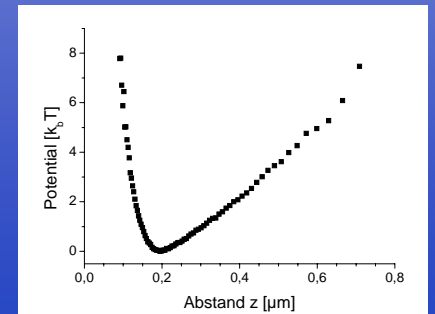


Histogram



Inversion of the Boltzman distribution

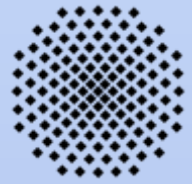
Potential



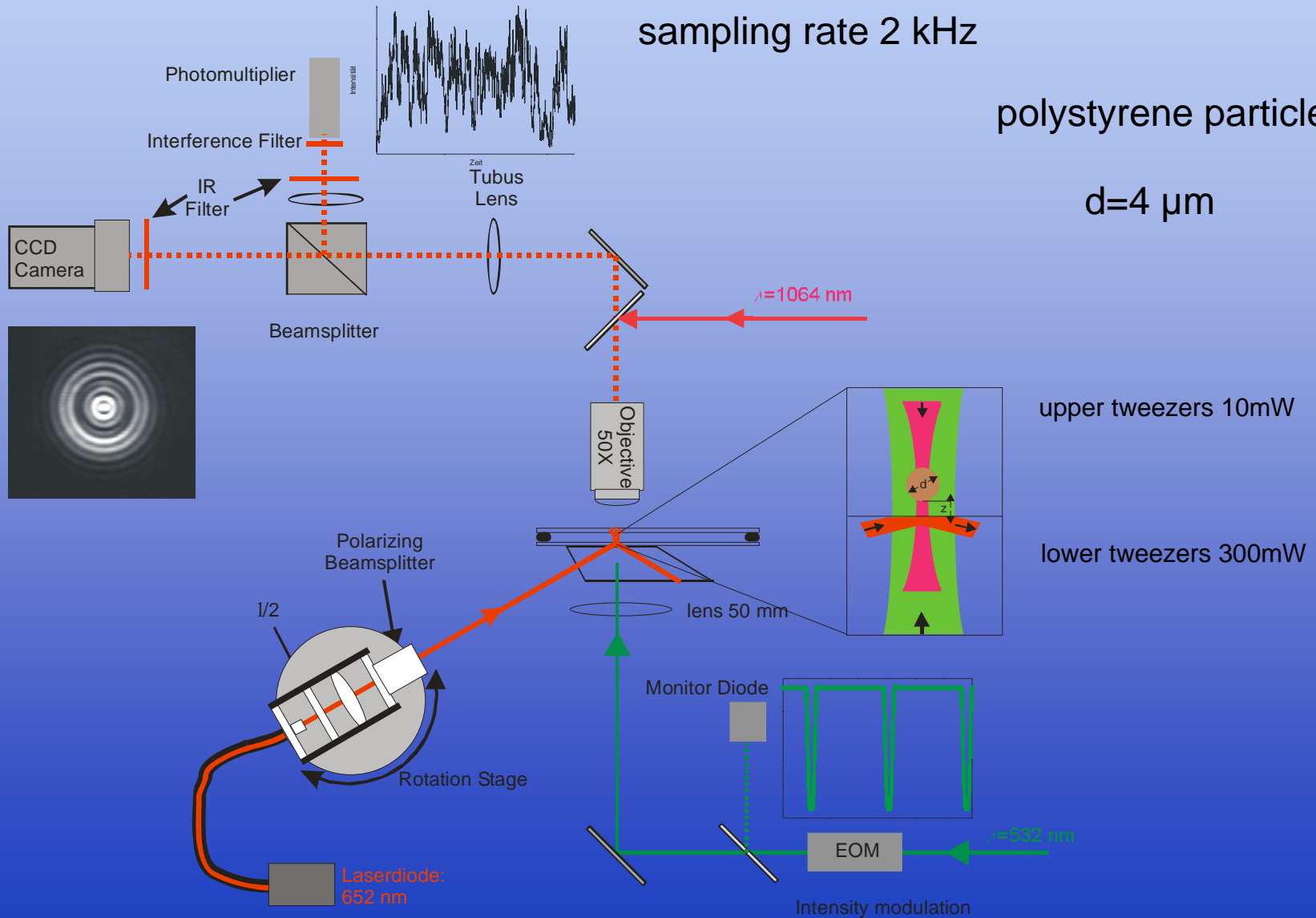
- Creation of an evanescent wave
- Exponentially decaying scattering intensity

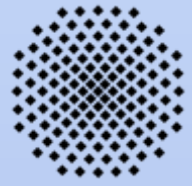


- Determination of the particle wall distance  $z$

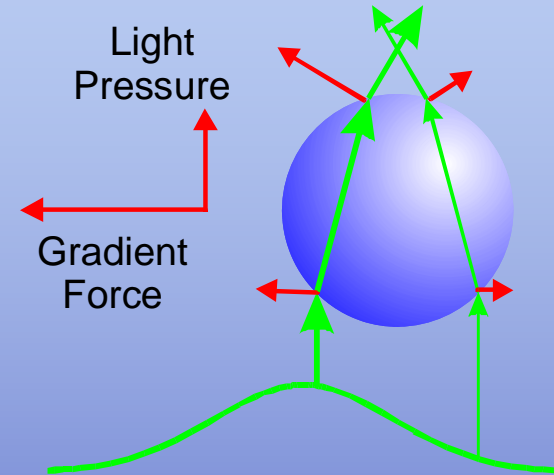
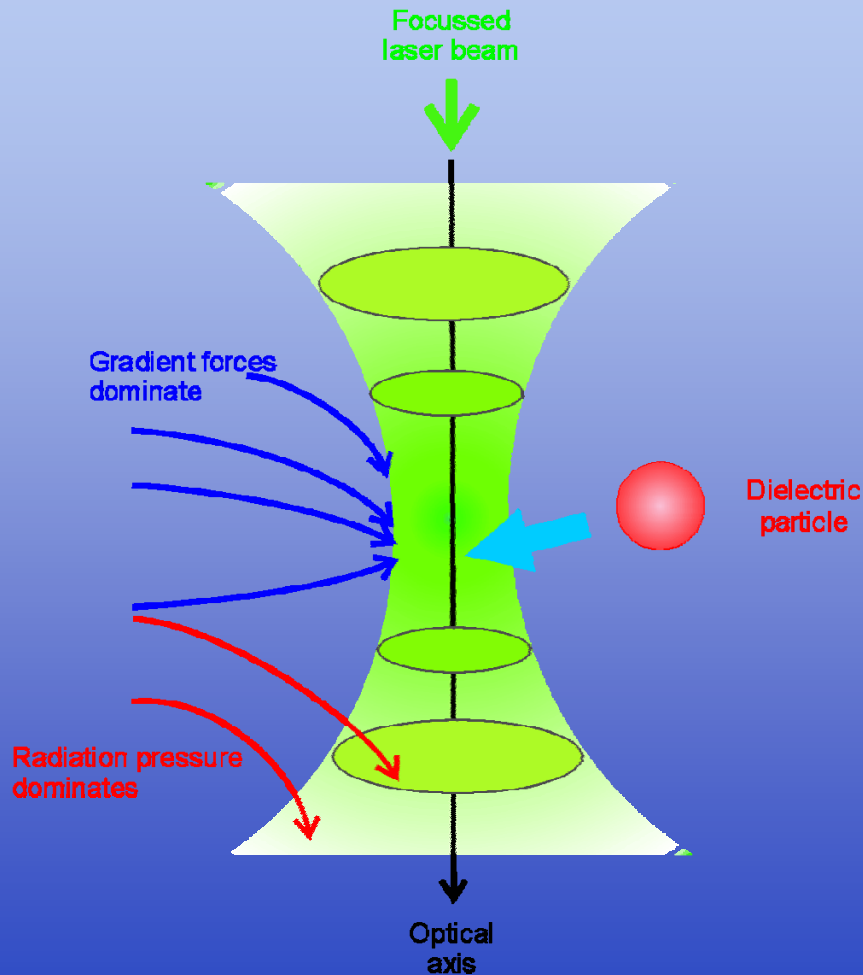


# experiment





# optical tweezers

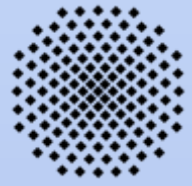


Dipole interaction energy for a particle in an electric field:

$$W \propto -\alpha \int E^2 dV$$

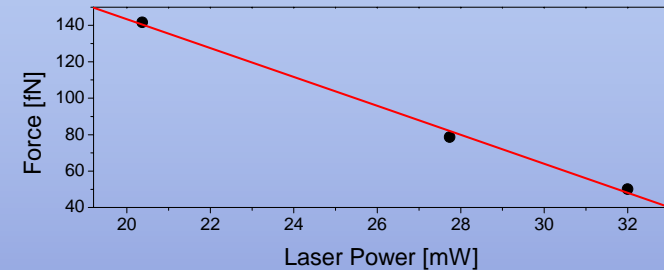
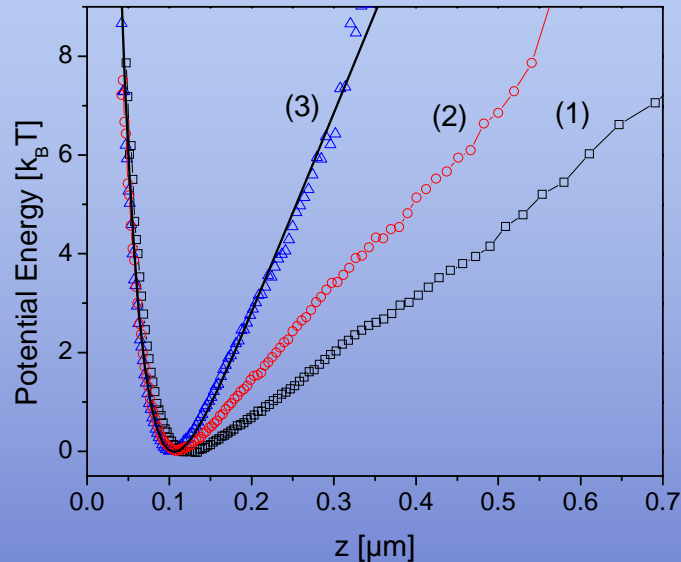
$$\alpha = \epsilon_{\text{particle}} / \epsilon_{\text{outside}} - 1$$

⇒ Particle moves towards higher fields.



# lower tweezers calibration

- increasing tweezers intensity

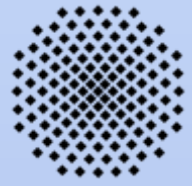


- linear dependance of light pressure
- tweezers calibration

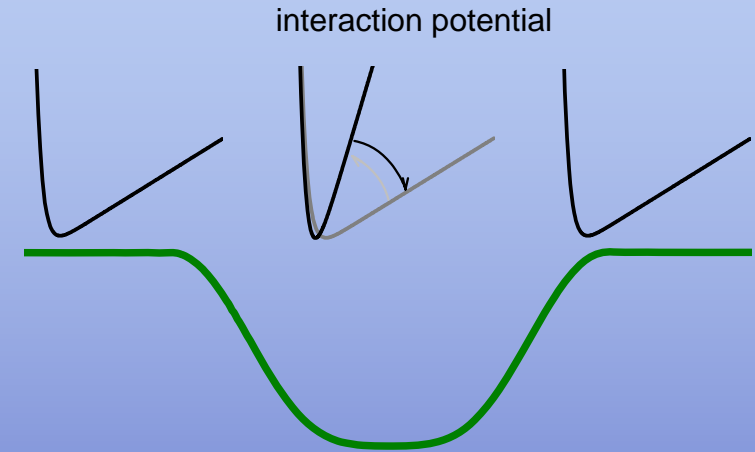
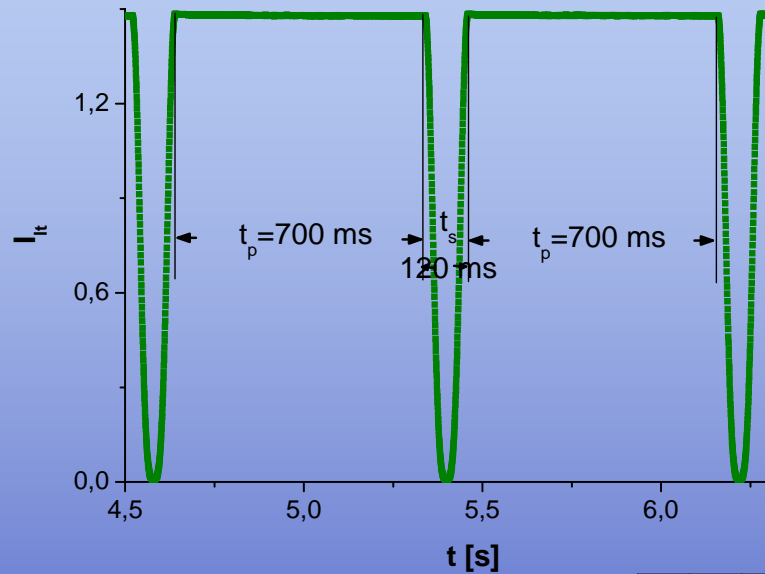
$$V_l = c \cdot I_l \cdot z$$

interaction potential

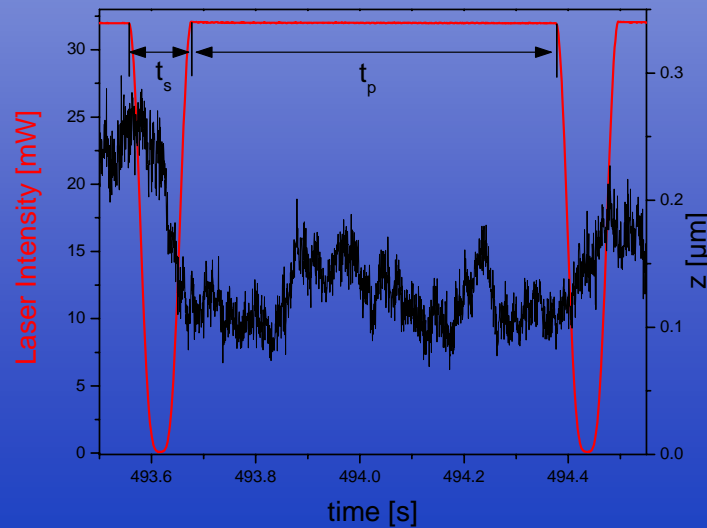
$$V[z(\tau), \lambda(\tau)] = A \cdot e^{-\kappa z} + B_0 \cdot z + I_{lt}[\lambda(\tau)] \cdot c \cdot z$$

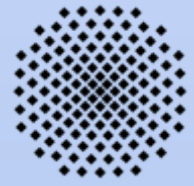


# pulse protocol



- strong coupling to the bath



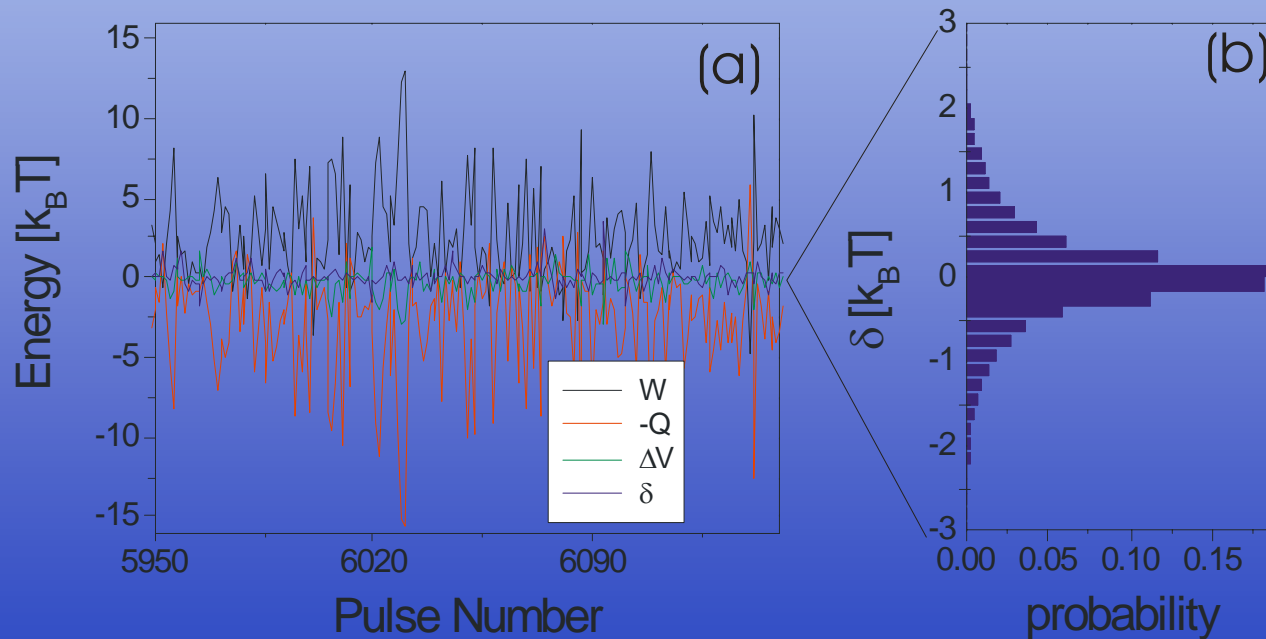


# energy conservation

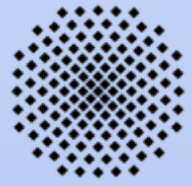
$$W = - \int_{t_0}^{t_e} \frac{\partial V}{\partial \lambda} \dot{\lambda}(t) d\tau = - \frac{c}{v} \sum_{\tau} \dot{I}_{lt} \cdot z(\tau)$$

$$Q = - \int_{t_0}^{t_e} \frac{\partial V}{\partial z} \dot{z} d\tau = - \frac{1}{v} \sum_{\tau} \frac{\partial V}{\partial z} \dot{z}(\tau)$$

energy balance:  $\delta = W - Q + \Delta V = 0$



energy resolution  
about 0.25 kT



# work distribution

particle wall potential is not harmonic

harmonic potentials



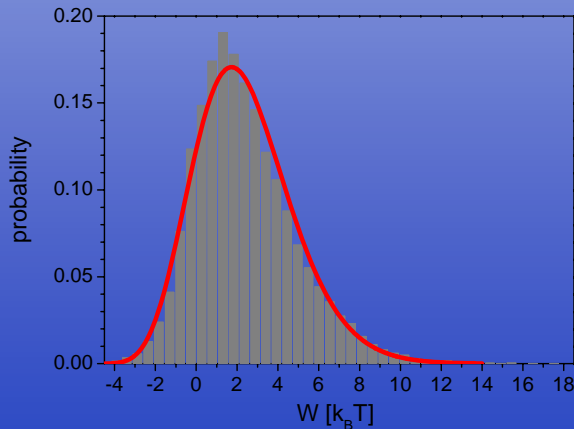
symmetric gaussian work distribution

- demonstrated with colloidal particles in 3d traps

non harmonic potentials



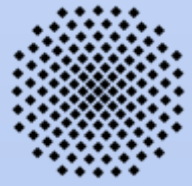
theory predicts:  
asymmetric non gaussian distribution



theory:

- no fit parameters

$$\langle W \rangle = 2.4 \text{ kT}$$



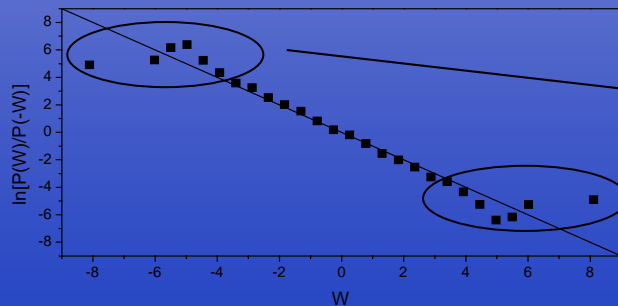
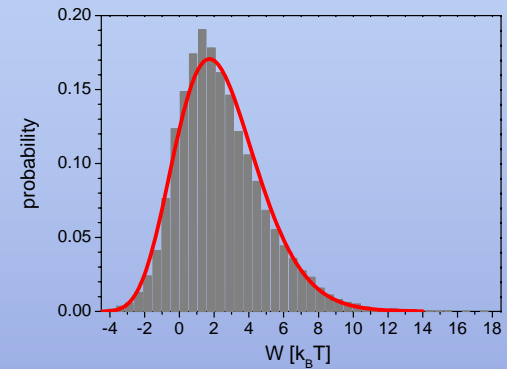
# fluctuation theorems

symmetric protocol:  $\Delta F = 0$

Jarzynski relation:

$$\text{exp. } \langle e^{-\beta W} \rangle = 1.03$$

detailed fluctuation theorem: 
$$\frac{P(-W)}{P(W)} = e^{-\beta W}$$



statistics!